# NECESSARY WEIERSTRASS CONDITIONS FOR ELLIPTIC SYSTEMS $\dagger$ 

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#### Abstract

Optimal control problems for objects described by systems of elliptic differential equations are considered. The controls are coefficients of leading terms in the equations and are constrained by equalities and inequalities. This formulation covers the optimal design problem for mechanical systems consisting of two materials with given volumes. The objective function is a functional which depends on the equations of state and their first derivatives. The main result is the construction of necessary Weierstrass conditions on inclusions of domains of one material in a domain of the other, which enables one to construct an approximation to the optimal solution. The case if a single $n$-dimensional elliptic equation is investigated in detail, for which a necessary Weierstrass condition is obtained when the inclusion is an $n$-dimensional ellipsoid. The application of a necessary Weierstrass condition to two optimal design problems is considered: the minimization of the work of external influences and the maximization of the torsional rigidity of a prismatic rod.


The traditional approach to the solution of optimal design problems consists of extending the set of admissible controls $\ddagger$ within which the optimal solution exists. In most cases such a solution is of purely theoretical interest, and can be used to estimate the optimal value of the functional. Attempting to approximate it by using some kind of regular solution is impracticable because in the approximation one has to deal with a problem of high dimension.

This paper focuses on the problem of obtaining a necessary Weierstrass condition for a strong minimum, which, together with the necessary conditions for a weak minimum, can be used to obtain a progressive improvement in the connectedness of the domains of inclusion of one material in the other [1].

## 1. STATEMENT OF THE PROBLEM

Let $\mathbb{R}^{n}$ and $\mathbb{R}^{m}$ be $n$ - and $m$-dimensional Cartesian vector spaces, respectively, i.e. real vector spaces of ordered sets of real numbers $x=\left(x_{1}, \ldots, x_{n}\right)$ and $u=\left(u_{1}, \ldots, u_{m}\right)$. Scalar products and norms in these spaces are defined in the usual manner

$$
(x, y)=x_{i} y_{i},|x|=(x, x)^{1 / 2},(u, v)=u_{i} v_{i},|u|=(u, u)^{1 / 2}
$$

Here and below the subscripts $i$ and $j$ take values from 1 to $n$, and the subscripts $k, l$ take values from 1 to $m$. Repeated $i$ or $j$ indices in a product imply summation from 1 to $n$, and repeated $k$ or $l$ subscripts imply summation from 1 to $m$.

We denote by $\Omega \subset \mathbb{R}^{n}$ a regular domain with piecewise-smooth boundary $\Gamma$ [2]. We assume that a vector function $f=\left(f_{1}(x), \ldots, f_{m}(x)\right) \in L^{2}(\Omega)$ is specified on $\Omega$, together with a vector-function $F=$ $\left(F_{1}(x), \ldots, F_{m}(x)\right) \in L^{2}\left(\Gamma_{F}\right)$ on $\Gamma_{F} \subset \Gamma$, where $L^{2}(\Omega), L^{2}\left(\Gamma_{F}\right)$ are Hilbert spaces of vector functions with scalar products and norms

$$
\begin{aligned}
& (f, g)=\int_{\Omega} f_{k} g_{k} d x,\|f\|_{\Omega}=\left(\int_{\Omega}|f|^{2} d x\right)^{1 / 2} \\
& (F, G)=\int_{\Gamma_{F}} F_{k} G_{k} d x,\|F\|_{\Gamma_{F}}=\left(\int_{\Gamma_{F}}|F|^{2} d \Gamma\right)^{1 / 2}
\end{aligned}
$$

$\dagger$ Prikl. Mat. Mekh. Vol. 59, No. 5, pp. 742-749, 1995.
$\ddagger$ LUR'YE K. A. and CHERKEYEV A. V., Composition problem for an optimal isotropic multiphase composite. Preprint No. 895, Leningrad. Phys.-Techn. Inst., Academy of Sciences of the U.S.S.R., 1984.

In the domain $\Omega$ we consider vector-functions $u=\left(u_{1}(x), \ldots, u_{m}(x)\right) \in H^{1}(\Omega)$ where $H^{1}(\Omega)$ is the Sobolev space of vector functions with scalar product and norm

$$
(u, v)=\int_{\Omega}\left(u_{k} v_{k}+u_{k, i} v_{k, i}\right) d x,\|u\|=(u, u)^{1 / 2}
$$

containing the subspace

$$
V(\Omega)=\left\{u \in H^{1}(\Omega) \mid u(x)=0, x \in \Gamma_{u}\right\}
$$

where $\Gamma_{u} \subset \Gamma, \Gamma_{u} \cap \Gamma_{F}=\varnothing$ and mes $\Gamma_{u}>0$.
In $V(\Omega)$ we define the symmetric bilinear form $A(u, v)=a_{i j k} u_{k i,} v_{i j}$ in which the coefficients $a_{i j k l}$ possess the symmetry properties $a_{i j k l}=a_{j k i l}=a_{i j k}=a_{j i k}$, and which satisfies the ellipticity condition

$$
\begin{equation*}
A(w, w) \geqslant \alpha\|w\|^{2}, \alpha>0 \tag{1.1}
\end{equation*}
$$

We define $u \in V(\Omega)$ to be the solution of the integral identity

$$
\begin{equation*}
\int_{\Omega}[A(u, w)-(f, w)] d x-\int_{\Gamma_{F}}(F, w) d \Gamma=0, \quad \forall w \in V(\Omega) \tag{1.2}
\end{equation*}
$$

If $A(u, w)$ satisfies condition (1.1), and $f \in L^{2}(\Omega), F \in L^{2}\left(\Gamma_{F}\right)$, then a unique solution $u \in V(\Omega)$ of the integral identity (1.2) exists [3].

We will now consider the formulation of the optimal design problem.
Suppose that the coefficients $a_{i k i l}^{(1)}$ and $a_{i k i l}^{(2)}$ are specified in measurable subspaces $\Omega_{1} \subset \Omega$ and $\Omega_{2} \subset$ $\Omega$ which satisfy the conditions $\Omega_{1} \cap \Omega_{2}=\varnothing, \bar{\Omega}_{1} \cup \bar{\Omega}_{2}=\bar{\Omega}$ and

$$
\begin{equation*}
\operatorname{mes} \Omega_{1}=\lambda_{1}, \operatorname{mes} \Omega_{2}=\lambda_{2}=\operatorname{mes} \Omega-\lambda_{1} \tag{1.3}
\end{equation*}
$$

It is required to solve the problem

$$
\begin{equation*}
\left.\inf . J(u), J(u)=\int_{\Omega_{1}} \varphi\left(u, \sigma^{(1)}(u)\right) d x+\int_{\Omega_{2}} \varphi(u)\right) d x+\int_{\Gamma_{F}} \varphi(u) d \Gamma \tag{1.4}
\end{equation*}
$$

where $u$ satisfies the integral identity

$$
\begin{align*}
& \int_{\Omega_{1}} A_{1}(u, w) d x+\int_{\Omega_{2}} A_{2}(u, w) d x-\int(f, w) d x-\int_{\Omega}(F, w) d \Gamma=0  \tag{1.5}\\
& \forall w \in V(\Omega)
\end{align*}
$$

which is obtained from (1.2) and in which $A(u, w)=a_{i j j}^{s} \mu_{k ; j} w_{l j}$. The functions $\varphi$ and $\psi$ in (1.4) are assumed to be differentiable $n$ times with respect to all their arguments, and $\sigma^{(s)}(u)$ is the matrix with components $\sigma_{i k}^{(s)}=a_{l k j l}^{(s)} \mu_{l, j}(s=1,2)$.

The sets $\Omega_{1}$ and $\Omega_{2}$ can have a fairly complicated structure and in certain cases can be found using averaging methods (cf. the publication cited in the footnote). The core of the analysis performed below is the determination of the sensitivity of functional (1.4) to inclusions of domains with coefficients $a_{i k j l}^{(2)}$ in the domain $\Omega_{1}$ and domains with coefficients $a_{i k j}^{(1)}$ in the domain $\Omega_{2}$.

We shall now assume that in the domain $\Omega$ there are two regular domains $\Omega_{1}$ and $\Omega_{2}$ filled with the first and second materials respectively, that $\Gamma_{12}$ is the boundary which separates the domains $\Omega_{1}$ and $\Omega_{2}$, and that $r=\left(r_{1}, \ldots, r_{n}\right)$ is the unit vector normal to it pointing from domain $\Omega_{1}$ into domain $\Omega_{2}$. We shall further assume that in the domains $\Omega_{1}$ and $\Omega_{2}$ the solution $u$ to the integral identity (1.5) is a twice-differentiable vector-function with respect to the arguments $x_{i}$. Then, using the identity

$$
a_{i k j l}^{(s)} u_{k, i} w_{l, j}=\left(a_{i k j}^{(s)} u_{k, i} w_{l}\right)_{, j}-a_{i k j l}^{(s)} u_{k, i j} w_{l}
$$

we obtain from (1.5) an elliptic boundary-value problem in the form of a system of differential equations with boundary conditions

$$
a_{i k j}^{(s)} u_{k, j j}^{(s)}+f_{\mathrm{l}}=0, x \in \Omega_{s}, s=1,2
$$

$$
\begin{gathered}
u_{k}^{(s)}=0, \quad x \in \Gamma_{u}, r_{j} \sigma_{j k}^{(s)}=F_{k}, \quad x \in \Gamma_{F} \\
r_{j} \sigma_{j k}^{(s)}=0, \quad x \in \Gamma_{s 0}=\Gamma_{s} \backslash\left(\bar{\Gamma}_{12} \cup \bar{\Gamma}_{F} \cup \bar{\Gamma}_{u}\right) \\
u_{k}^{(1)}=u_{k}^{(2)}, \quad r_{j} \sigma_{j k}^{(1)}=r_{j} \sigma_{j k}^{(2)}, x \in \Gamma_{12} ; \Gamma_{12}=\Gamma_{1} \cap \Gamma_{2}, u=u^{(s)}, x \in \Omega_{s}
\end{gathered}
$$

where $\Gamma_{1}$ and $\Gamma_{2}$ are respectively the boundaries of the domains $\Omega_{1}$ and $\Omega_{2}$.

## 2. NECESSARY OPTIMALITY CONDITIONS

We shall assume that there is an optimal solution such that the two domains $\Omega_{1}^{*}$ and $\Omega_{2}^{*}$ with common boundary $\Gamma_{12}^{*}\left(\Gamma_{12}^{*}=\Gamma_{1}^{*} \cap \Gamma_{2}^{*}\right)$ are filled with the first and second materials, respectively.

We construct an extended functional $I(u)$ in which the left-hand side of the identity (1.5) has been subtracted from the functional $J(u)$, and compute its first variation

$$
\begin{align*}
& \delta I=\chi_{1}(w, \delta u)+\chi_{2}(w, \delta u)+\int_{\Gamma_{F}} \frac{\partial \psi}{\partial u_{k}} \delta u_{k} d \Gamma+\int_{\Gamma_{i}^{*} \cap \Gamma}\left[\Lambda_{1}\left(u^{*}, w\right)-\Lambda_{2}\left(u^{*}, w\right)\right] \delta r d \Gamma+ \\
& +\int_{\Gamma_{1}^{*} \cap \Gamma}\left[\Lambda_{2}\left(u^{*}, w\right)-\Lambda_{1}\left(u^{*}, w\right)\right] \delta r d \Gamma+\int_{\Gamma_{i 2}^{*}}\left[\Lambda_{1}\left(u^{*}, w\right)-\Lambda_{2}\left(u^{*}, w\right)\right] \delta r d \Gamma  \tag{2.1}\\
& \chi_{s}(w, \delta u)=\int_{\Omega_{i}^{*}}\left[\frac{\partial \varphi}{\partial u_{k}}\left(u^{*}, \sigma^{(s)}\left(u^{*}\right)\right) \delta u_{k}+\frac{\partial \varphi}{\partial \sigma_{i k}^{(s)}}\left(u^{*}, \sigma^{(s)}\left(u^{*}\right)\right) \sigma_{i k}^{(s)}(\delta u)+A_{s}(w, \delta u)\right] d x, \\
& \Lambda_{s}\left(u^{*}, w\right)=\varphi\left(u^{*}, \delta^{(s)}\left(u^{*}\right)\right)-A_{s}\left(u^{*}, w\right), s=1,2
\end{align*}
$$

Here $u^{*}$ is the optimal solution of the problem, $\delta u$ is the variation of $u^{*}$, and $\delta r$ is a variation of the boundary $\Gamma_{12}^{*}$ satisfying the condition [1]

$$
\begin{equation*}
\int_{\Gamma_{12}^{*}} \delta r d \Gamma+\int_{\Gamma_{i} \cap \Gamma} \delta r d \Gamma+\int_{\Gamma_{2}^{*} \cap \Gamma} \delta r d \Gamma=0 \tag{2.2}
\end{equation*}
$$

which follows from condition (1.3).
In (2.1) we put $w=v^{*}$, where $v^{*}$ satisfies the integral identity

$$
\begin{equation*}
\chi_{1}\left(v^{*}, w\right)+\chi_{2}\left(v^{*}, w\right)+\int_{\Gamma_{F}} \frac{\partial \psi}{\partial u_{k}}\left(u^{*}\right) w_{k} d \Gamma=0, \quad \forall w \in V(\Omega) \tag{2.3}
\end{equation*}
$$

Then, using $\delta u \in V(\Omega)$, we obtain the inequality

$$
\begin{gather*}
\delta I=\int_{\Gamma_{1}^{*} \cap \Gamma}\left[\Lambda_{1}\left(u^{*}, v^{*}\right]-\Lambda_{2}\left(u^{*}, v^{*}\right] \delta r d \Gamma+\int_{\Gamma_{2}^{*} \cap \Gamma}\left[\Lambda_{1}\left(u^{*}, v^{*}\right]-\Lambda_{2}\left(u^{*}, v^{*}\right] \delta r d \Gamma+\right.\right. \\
+\int_{\Gamma_{12}^{*}}\left[\Lambda_{1}\left(u^{*}, v^{*}\right]-\Lambda_{2}\left(u^{*}, v^{*}\right] \delta r d \Gamma \geqslant 0\right. \tag{2.4}
\end{gather*}
$$

which must be satisfied by any variations of the boundary $\Gamma_{12}^{*}, \Gamma_{1}^{*} \cap \Gamma, \Gamma_{2}^{*} \cap \Gamma$ satisfying equality (2.2). It has been shown [2] that the conditions

$$
\begin{array}{ll}
\Lambda_{1}\left(u^{*}, v^{*}\right)-\Lambda_{2}\left(u^{*}, v^{*}\right) \leqslant \zeta, & x \in \Gamma_{1}^{*} \cap \Gamma \\
\Lambda_{1}\left(u^{*}, v^{*}\right)-\Lambda_{2}\left(u^{*}, v^{*}\right) \geqslant \zeta, & x \in \Gamma_{2}^{*} \cap \Gamma \\
\Lambda_{1}\left(u^{*}, v^{*}\right)-\Lambda_{2}\left(u^{*}, v^{*}\right)=\zeta, & x \in \Gamma_{12}^{*} \tag{2.5}
\end{array}
$$

are necessary and sufficient for (2.2) and (2.4) to be satisfied.

## 3. NECESSARY WEIERSTRASS CONDITIONS

We will analyse a variation of the functional in which a small domain of the first material is included
in the domain of the second material, and vice versa.
We take $\forall x_{0} \in \Omega_{1}^{*}\left(\forall x_{0} \in \Omega_{2}^{*}\right)$ and a convex domain $\Omega_{0}(\eta)$ all of whose geometrical dimensions can vary in proportion to $\eta$. Then

$$
\begin{equation*}
\text { mes } \Omega_{0}(\eta)=\eta^{n} \text { mes } \Omega_{0}\left(\Omega_{0}=\Omega_{0}(1)\right) \tag{3.1}
\end{equation*}
$$

In order for condition (1.3) to remain satisfied, it is necessary to vary the boundary $\Gamma_{12}^{*}$ by an amount $r(y, \eta), y \in \Gamma_{12}^{*}[2]$, where $r(y, \eta)<0$ in the case of an inclusion of the first material in the second, and $r(y, \eta)>0$ in the case of an inclusion of the second material in the first.

In this case functional (1.4) can be reduced to the form

$$
\begin{align*}
& I=\int_{\Gamma_{F}}[\psi(u)-(F, w)] d \Gamma-\int_{\Omega}(f, w) d x+\int_{\Omega_{1}} \Lambda_{1}(u, w) d x+ \\
& +\int_{\Omega_{2}} \Lambda_{2}(u, w) d x \pm \int_{\Omega_{0}(\eta)}\left[\Lambda_{1}(u, w)-\Lambda_{2}(u, w)\right] d x \tag{3.2}
\end{align*}
$$

The plus sign in front of the last integral corresponds to the inclusion of the first material in the second, and the minus sign to the inclusion of the second material in the first.

In order to satisfy condition (1.3) the function $r(y, \eta)$ must be proportional to $\eta^{n}$, hence $\delta r=$ .. $\delta^{n-1} r=0$. The variations $\delta u, \ldots, \delta^{n} u$ are given by integral identities in which the external actions are functions proportional to $\delta r, \ldots, \delta^{n} r$, and hence $\delta u=\ldots=\delta^{n-1} u=0[1]$.

Substituting $w=v^{*}$ into functional (3.2) and sequentially computing the variations $\delta I, \ldots, \delta^{n} I$, we obtain

$$
\begin{aligned}
& \delta l=\ldots=\delta^{n-1} l=0 \\
& \delta^{n} I=\int_{\Gamma_{i 2}^{*}}\left[\Lambda_{1}\left(u^{*}, v^{*}\right)-\Lambda_{2}\left(u^{*}, v^{*}\right)\right] \delta^{n} r d \Gamma \pm \frac{d^{n}}{d \eta^{n}}\left(\int_{\Omega_{1}(\eta)}\left[\Lambda_{1}\left(u^{*}, v^{*}\right)-\Lambda_{2}\left(u^{*}, v^{*}\right)\right] d x\right)_{\eta=0}
\end{aligned}
$$

Using the necessary conditions $\delta^{n} I \geqslant 0$ and (2.5) we find

$$
\begin{equation*}
\pm \frac{d^{n}}{d \eta^{n}} \int_{\Omega_{0}(\eta)}\left[\Lambda_{1}\left(u^{*}, v^{*}\right)-\Lambda_{2}\left(u^{*}, v^{*}\right)\right] d x \geqslant \pm n!\zeta \operatorname{mes} \Omega_{0} \tag{3.3}
\end{equation*}
$$

The minus sign is taken for points $x_{0} \in \Omega_{1}^{*}$ and the plus sign for $x_{0} \in \Omega_{2}^{*}$.
In the left-hand side of inequality (3.3) the integral is evaluated over the domain $\Omega_{0}(\eta)$ (an $n$-dimensional volume proportional to $\eta^{n}$ ) of an expression which depends on the vector function $u$, which is a perturbation of the solution to the integral identity (1.5) that has been produced by the inclusion $\Omega_{0}(\eta)$.

We will obtain integral identities for determining $u$. To do this we consider an inclusion $\Omega_{0}(\eta)$ of the second material in the domain $\Omega_{1}^{*}$. Subtracting from the integral identity for this case the integral identity for the optimal solution $u^{*}$ and transforming it, we find

$$
\begin{align*}
& \Omega_{1}^{*} \backslash \frac{\Omega}{\Omega}_{0(\eta)} A_{1}\left(u^{0}-u^{*}, w\right) d x+\int_{\Omega_{0}}\left(u^{0}-u^{*}, w\right) d x=\int_{\partial \Omega_{0}} r_{i}\left[\sigma_{i l}^{(1)}\left(u^{*}\left(x_{0}\right)\right)-\sigma_{i l}^{(2)}\left(u^{*}\left(x_{0}\right)\right)\right] w_{l} d \Gamma,  \tag{3.4}\\
& \forall x_{0} \in \Omega_{1}^{*}
\end{align*}
$$

It is of course impossible to determine $u$ exactly in the integral identity (3.4). However, when $\eta \rightarrow 0$ we have $u \rightarrow u^{0}$ for which the integral identity

$$
\begin{align*}
& \int_{\mathbb{R}^{n} \backslash \Omega_{0}} A_{1}\left(u^{0}-u^{*}, \omega\right) d x+\int_{\Omega_{0}} A_{2}\left(u^{0}-u^{*}, \omega\right) d x=\int_{\partial \Omega_{i}} r_{i}\left[\sigma_{i l}^{(1)}\left(u^{*}\left(x_{0}\right)\right)-\sigma_{i l}^{(2)}\left(u^{*}\left(x_{0}\right)\right)\right] \omega_{i} d \Gamma, \\
& \forall x_{0} \in \Omega_{1}^{*} \tag{3.5}
\end{align*}
$$

holds, where the $r_{i}$ are the components of the external unit normal to the boundary $\partial \Omega_{0}$ of the domain $\Omega_{0}$.

Performing similar constructions for an inclusion $\Omega_{0}(\eta)$ of the first material into the second, we obtain an integral identity differing from (3.5) by the replacement of the superscript 1 by the superscript 2 and vice versa.

Using the latter two identities together with inequality (3.3), we finally obtain the inequalities

$$
\begin{align*}
& \Lambda_{1}\left(u^{0}, v^{0}\right)-\Lambda_{2}\left(u^{0}, v^{*}\right) \leqslant \zeta, x \in \Omega_{1}^{*} \\
& \Lambda_{1}\left(u^{0}, v^{0}\right)-\Lambda_{2}\left(u^{0}, v^{*}\right) \geqslant \zeta, x \in \Omega_{2}^{*}  \tag{3.6}\\
& u^{*}=v^{*}\left(x_{0}\right), v^{*}=v^{*}\left(x_{0}\right), u^{0}=u^{0}\left(x_{0}\right)
\end{align*}
$$

Inequalities (3.6) are called necessary Weierstrass conditions. The first must be satisfied by any $x_{0} \in$ $\Omega_{1}^{*}$ and the second for any $x_{0} \in \Omega_{2}^{*}$.

## 4. NECESSARY WEIERSTRASS CONDITIONS FOR AN ELLIPSOIDAL INCLUSION ( $m=1$ )

For the case when $m=1$ the bilinear form $A(u, v)=a_{i j} u_{j} v_{, j}$ can be reduced to the form

$$
\begin{equation*}
A(u, v)=h u_{i} v_{i} \tag{4.1}
\end{equation*}
$$

by a change of coordinates.
Below we shall assume that the Cartesian system of coordinates $x_{i}$ has been chosen from the start so that the bilinear form $A(u, v)$ has the form (4.1), and also $h_{2}>h_{1}$. The components of the vectors $\sigma^{(s)}$ are given by the relations $\sigma_{i}^{(1)}(u)=h_{1} u_{i}, \sigma_{i}^{(2)}(u)=h_{2} u_{i}$ and the necessary Weierstrass conditions (3.6) are the inequalities

$$
\begin{align*}
& \Delta h u_{i, i}^{0} v_{, i}^{*}-\varphi\left(u^{*}, \sigma^{(2)}\left(u^{0}\right)\right)+\varphi\left(u^{*}, \sigma^{(1)}\left(u^{0}\right)\right) \leqslant \zeta \\
& \Delta h u_{, i}^{0} v_{, i}^{*}-\varphi\left(u^{*}, \sigma^{(2)}\left(u^{0}\right)\right)+\varphi\left(u^{*}, \sigma^{(1)}\left(u^{0}\right)\right) \geqslant \zeta \tag{4.2}
\end{align*}
$$

where $\Delta h=h_{2}-h_{1}$.
In this case the solution $u^{0}$ of the integral identities (3.5) and (3.6) are known [4]

$$
\begin{align*}
& u^{0}=u^{*}+\kappa_{i}^{(s)} x_{i} u_{i}^{*}\left(x_{0}\right), s=1,2 \\
& \kappa_{i}^{(1)}=-\Delta h \mu_{i}\left(h_{1}+\Delta h \mu_{i}\right)^{-1}, \kappa_{i}^{(2)}=\Delta h \mu_{i}\left(h_{2}+\Delta h \mu_{i}\right)^{-1}  \tag{4.3}\\
& \mu_{i}=\frac{1}{2} a_{1} \ldots a_{n} \int_{0}^{\infty} \frac{d \rho}{\left(a_{i}^{2}+\rho\right) \sqrt{\left(a_{1}^{2}+\rho\right) \ldots\left(a_{n}^{2}+\rho\right)}}
\end{align*}
$$

where the superscript 1 in parentheses corresponds to an inclusion of the second material in the first and the 2 corresponds to the first material in the second.

Without loss of generality we can assume that the semi-axes of the ellipsoid have been ordered as follows:

$$
\begin{equation*}
a_{1}=1, a_{1} \geqslant a_{2} \geqslant \ldots a_{n} \geqslant 0 \tag{4.4}
\end{equation*}
$$

It then follows from (4.3) that $0 \leqslant \mu_{1} \leqslant \mu_{2} \leqslant \ldots \mu_{n} \leqslant 1$, the last inequality being obtained from an estimate of the integrand in integral (4.3)

$$
\mu_{n} \leqslant \frac{a_{n}}{2} \int_{0}^{\infty} \frac{d \rho}{\left(a_{n}^{2}+\rho\right)^{3 / 2}}=1
$$

Analysis shows that $\kappa^{(1)}(\mu)$ is a monotonically decreasing, and $\kappa^{(2)}(\mu)$ is a monotonically increasing function of $\mu$ when $0 \leqslant \mu \leqslant 1$ (Fig. 1).
We will consider the case when function $\varphi=\varphi(u)$. Then inequalities (4.2) take the respective forms

$$
\begin{align*}
& \Delta h u_{, i}^{*} v_{i,}^{*} \leqslant \zeta \pm \Delta h \kappa_{i}^{(1)} u_{i,}^{*} v_{, i}^{*}, \quad \forall x_{0} \in \Omega_{1}^{*} \\
& \Delta h u_{, i}^{*} v_{i,}^{*} \geqslant \zeta \pm \Delta \kappa_{i}^{(2)} u_{i,}^{*} v_{i,}^{*}, \quad \forall x_{0} \in \Omega_{2}^{*} \tag{4.5}
\end{align*}
$$



Fig. 1.


Fig. 2.
(where $u^{*}=u^{*}\left(x_{0}\right), v^{*}=v^{*}\left(x_{0}\right)$ ). We denote by $\alpha_{i}$ the angles between $\nabla u^{*}\left(x_{0}\right)$-the gradient of the function $u^{*}$ at the point $x_{0}$-and the axes of the $n$-dimensional ellipsoid, by $\beta_{i}$ the angles between $\nabla u^{*}\left(x_{0}\right)$ -the gradient of the function $v^{*}$ at the point $x_{0}$-and the axes of the $n$-dimensional ellipsoid, and by $\gamma$ the angle between the gradients $\nabla u^{*}\left(x_{0}\right)$ and $\nabla u^{*}\left(x_{0}\right)$. Then inequalities (4.5) can be represented in the form

$$
\begin{align*}
& \Delta h\left|\nabla u^{*} \| \nabla v^{*}\right| \cos \gamma \leqslant \zeta-\Phi_{1}\left(\alpha_{i}, \beta_{i}, a_{i}\right) \\
& \Delta h\left|\nabla u^{*} \| \nabla v^{*}\right| \cos \gamma \leqslant \zeta-\Phi_{2}\left(\alpha_{i}, \beta_{i}, a_{i}\right)  \tag{4.6}\\
& \Phi_{s}=\Delta h\left|\nabla u^{*} \| \nabla v^{*}\right| \kappa_{i}^{(s)} \cos \alpha_{i} \cos \beta_{i}, s=1,2
\end{align*}
$$

Inequalities (4.6) must be satisfied by all $\alpha_{i}, \beta_{i}, a_{i}$ satisfying the equalities

$$
\begin{equation*}
\sum_{i=1}^{n} \cos ^{2} \alpha_{i}=1, \sum_{i=1}^{n} \cos ^{2} \beta_{i}=1, \sum_{i=1}^{n} \cos \alpha_{i} \cos \beta_{i}=\cos \gamma \tag{4.7}
\end{equation*}
$$

and inequalities (4.4).
We will solve the two problems

$$
\max _{\alpha_{i}, \beta_{i}, a_{i}} \Phi_{1}\left(\alpha_{i}, \beta_{i}, a_{i}\right), \min _{\alpha_{i}, \beta_{i}, a_{i}} \Phi_{2}\left(\alpha_{i}, \beta_{i}, a_{i}\right)
$$

for $\alpha_{i}, \beta_{i}, a_{i}$ satisfying (4.7) and (4.4). Then inequalities (4.6) hold for all $\alpha_{i}, \beta_{i}, a_{i}$ satisfying conditions (4.7) and (4.4). The maximum of the function $\Phi_{1}\left(\alpha_{i}, \beta_{i}, a_{i}\right)$ and the minimum of the function $\Phi_{2}\left(\alpha_{i}, \beta_{i}\right.$, $a_{i}$ ) are reached at the points (Fig. 2)

$$
\begin{align*}
& a_{1}=\ldots=a_{n-1}=1, a_{n}=0, \alpha_{1}=-\beta_{1}=\gamma / 2 \\
& \alpha_{2}=\beta_{2}=\ldots=\alpha_{n-1}=\beta_{n-1}=\pi / 2, \alpha_{n}=\pi / 2-\gamma / 2, \beta_{n}=\pi / 2+\gamma / 2 \tag{4.8}
\end{align*}
$$

and are given by

$$
\Phi_{1}^{*}=h_{2}^{-1} \Delta h^{2}\left|\nabla u^{*}\right|\left|\nabla v^{*}\right| \sin ^{2}(\gamma / 2), \Phi_{2}^{*}=-h_{1}^{-1} \Delta h^{2}\left|\nabla u^{*} \| \nabla v^{*}\right| \sin ^{2}(\gamma / 2)
$$

Analysis of relations (4.8) shows that the extremal values of the functions $\Phi_{1}$ and $\Phi_{2}$ are reached on an $n$-dimensional oblate spheroid the normal to which coincides with the $x_{n}$ axis. The necessary Weierstrass conditions finally take the form

$$
\begin{align*}
& \Delta h\left|\nabla u^{*}\right|\left|\nabla v^{*}\right| \cos \gamma \leqslant \zeta-h_{2}^{-1} \Delta h^{2}\left|\nabla u^{*}\right| \nabla v^{*} \mid \sin ^{2}(\gamma / 2)  \tag{4.9}\\
& \Delta h\left|\nabla u^{*}\right|\left|\nabla v^{*}\right| \cos \gamma \geqslant \zeta+h_{1}^{-1} \Delta h^{2}\left|\nabla u^{*}\right| \nabla v^{*} \mid \sin ^{2}(\gamma / 2)
\end{align*}
$$

where $u^{*}=u^{*}\left(x_{0}\right), v^{*}=v^{*}\left(x_{0}\right)$.
Example 1. Consider the problem of minimizing the work of $f$ and $F$ through a displacement $u$. In this case $\varphi$ $=f u, \psi=F u$. Analysis of integral identities (1.5) and (2.3) shows that $\gamma=0$, and the necessary Weierstrass conditions consequently have the simple form

$$
\begin{equation*}
\Delta h\left|\nabla u^{*}\right|^{2} \leqslant \zeta, x_{0} \in \Omega_{1}^{*} ; \Delta h\left|\nabla u^{*}\right|^{2} \geqslant \zeta, x_{0} \in \Omega_{2}^{*} \tag{4.10}
\end{equation*}
$$

We put $x_{0}=y_{-}=y-0 r$ in the first inequality and $x_{0}=y_{+}=y+0 r$ in the second, where $r$ is the unit normal to the boundary $\Gamma_{12}^{*}$, and $y \in \Gamma_{12}^{*}$. We also note that $\zeta=\left.h_{2}\left|\nabla u^{*}\right|^{2}\left(y_{+}\right)\right|^{2}-h_{1}\left|\nabla u^{*}\left(y_{-}\right)\right|^{2}$. Analysis of inequalities (4.10) shows that they can be satisfied if

$$
\partial u^{*}\left(y_{-}\right) / \partial r=\partial u^{*}\left(y_{+}\right) / \partial r=0, \forall y \in \Gamma_{12}
$$

Example 2. Consider the problem of maximizing the torsional rigidity of a rod composed of two materials, one flexible and one stiff, with shear moduli $G_{1}$ and $G_{2}$, respectively. In this case $\varphi=-2 u, f=2, h_{1}=\rho G_{1}, h_{2}=\rho G_{2}$, where the quantity $\rho$ is proportional to the torque applied to the rod. It follows from integral identities (1.5) and (2.3) that $v^{*}=u^{*}$. In this case $\gamma=\pi, \zeta=-\Delta\left|\nabla u^{*}(y)\right|^{2}$ for $y \in \Gamma_{11}^{*}$, and inequalities (4.9) acquire the form

$$
\begin{align*}
& h_{1} h_{2}^{-1}\left|\nabla u^{*}\left(x_{0}\right)\right|^{2} \geqslant \zeta, \quad \forall x_{0} \in \Omega_{1}^{*} \\
& h_{2} h_{1}^{-1}\left|\nabla u^{*}\left(x_{0}\right)\right|^{2} \leqslant \zeta, \forall x_{0} \in \Omega_{2}^{*} \tag{4.11}
\end{align*}
$$

We put $x_{0}=y_{-}=y-0 r$ in the first inequality and $x_{0}=y_{+}=y+0 r$ in the second, where $y \in \Gamma_{12}^{*}$ and $r$ is the unit normal to $\Gamma_{12}^{*}$. Then from (4.11) one can obtain the inequality

$$
\begin{aligned}
& h_{1}\left(2 h_{2}-h_{1}\right)\left|\nabla u^{*}\left(y_{-}\right)\right|^{2} \geqslant h_{2}^{2}\left|\nabla u^{*}\left(y_{+}\right)\right|^{2} \\
& h_{2}\left(2 h_{1}-h_{2}\right)\left|\nabla u^{*}\left(y_{+}\right)\right|^{2} \geqslant h_{1}^{2}\left|\nabla u^{*}\left(y_{-}\right)\right|^{2}
\end{aligned}
$$

from which it follows that they can only be satisfied in the case when

$$
u_{, i}^{*}\left(y_{-}\right)=u_{i,}^{*}\left(y_{+}\right)=0, i=1, \ldots, n
$$

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